## MODELLING AND SIMULATION OF STRAIN RESISTANCE OF ALLOYS TAKING INTO ACCOUNT BARRIER EFFECTS

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The paper proposes a model of strain resistance of alloy under high-temperature deformation. The model describes hardening of alloy due to the increase of dislocation density, as well as the barrier effect of blocking free dislocations, boundaries of grains and subgrains by dispersoids. The model also takes into account the softening processes associated with the recovery and dynamic recrystallization. The model has been tested on the rheological behavior of an Al-Mg alloy named AMg6 at temperatures of 400 and 500 °C in the range of strain rates from 5 to 25 s<sup>-1</sup>. It was found in this temperature – strain rate range that the curve of strain resistance of the AMg6 alloy consists of several portions. First there is hardening of the material, then there is material softening, which is again replaced by hardening of the material. With the use of the electron backscatter diffraction technique and transmission electron microscopy, it was found that the main process of softening at investigated temperatures is dynamic recrystallization. The appearance of the second portion of hardening on the strain resistance curve is the inhibition of dynamic recrystallization, as well as manifestation of the barrier effect of blocking free dislocations, grain and subgrain boundaries by dispersoids.

Keywords: Al-Mg alloy, AMg6, strain resistance model, strain resistance, rheology, high temperature deformation, recrystallization, microstructure, dispersoids, barrier effect.

#### 1. Introduction

Processes associated with material hardening and softening occur in metal materials under deformation at high temperatures. Depending on temperature – strain rate conditions of deformation, dynamic recovery, polygonization or recrystallization may prevail in a material, or these processes can take place simultaneously in different or equal portions [1–3]. The main mechanism of hardening in alloys is the process associated with an increase in the dislocation density and impeding the motion of free dislocations. The softening process can be inhibited due to the barrier effect connected with blocking free dislocation motion by impurity atoms and inclusions and with the growth of grains and subgrains by pinning on various inclusions [4]. The barrier effect of free dislocation blocking is usually associated with the Portevin–Le Chatelier and Savart–Masson effects and the inverse relation between strain resistance and strain rate [5–13]. The inhibition or blocking of grain and subgrain boundaries on inclusions generally occurs in metal matrix composites and high alloys [7, 11, 12, 14–24].

The physical processes of hardening and softening occurring under deformation change the form of the strain resistance curve and affect the final physical-mechanical properties of alloys [2, 25]. For example, the impeded growth of grains under deformation results in the prevalence of a certain grain orientation over the others thus improving the mechanical properties of products in this direction, and this may have a positive effect in the designing of crucial products. The barrier effect of dislocations blocked by inclusions or impurity atoms, as a rule, has a negative influence on the final technological properties, such as lower corrosion-resistance, strength and fatigue strength [13, 26–28]. Since the processes of structure formation affect the strain resistance curve, a mathematical model of strain resistance should include components describing the physical processes of hardening and softening. A large number of models have already been proposed in the literature,



which take into consideration separately the influence of discontinuous dynamic recrystallization [29–37], *in-situ* recrystallization [38–41] or the Portevin–Le Chatelier effect [5–7] on strain resistance. However, in the literature we have not found any models of strain resistance for alloys under high-temperature deformation that would describe the combined influence of hardening, recovery, recrystallization and barrier effects.

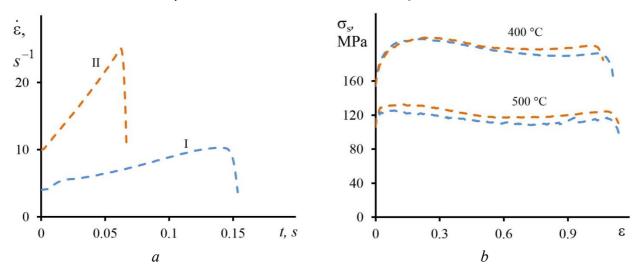
Previously, we developed and tested a model of strain resistance for a number of aluminum alloys and steels with an fcc lattice [42–44] within a wide temperature – strain rate range. The ability of the model to predict the part of the dynamic recrystallization with acceptable engineering precision under time-varying strain rate of the specimen was shown in [45].

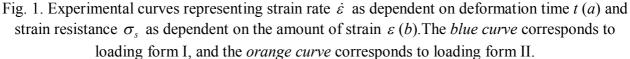
The paper proposes a new rheological model of strain resistance, which is based on the model found in [45] and takes into account the viscous properties of an alloy under plastic deformation, its softening due to recovery and recrystallization and its hardening due to the barrier effect of blocking the motion of free dislocations and the growth of recrystallized grains by inclusions. The model has been identified by experimental data obtained for the AMg6 alloy.

#### 2. Material and investigation procedure

Cylindrical specimens with a diameter of  $8 \pm 0.1$  mm and a height of  $12 \pm 0.1$  mm were machined from 10-mm AMg6 alloy rod (Al 92.4; Mg 6.56; Fe 0.27; Si 0.16; Cu 0.013; Mn 0.05; Ti 0.04; Zn 0.02; Be 0.0012%). They were compressed at 400 and 500 °C by a cam plastometer installed in the "Plastometriya" center for collective use at the Institute of Engineering Science, UB RAS. The specimens were annealed in vacuum at 500 °C for 4 hours before deformation.

To avoid heat losses during testing, the specimen was wrapped by kaolin wool and placed into a metal container together with heads. The container with the specimen inside was heated in an electric furnace up to the temperature of testing and then transferred from the furnace to the working space of the apparatus where the specimen was deformed. The specimen temperature during heating and the temperature variations inside the container during testing were monitored by a chromel-alumel thermocouple. Since we could not avoid specimen temperature fall while transferring the container from the furnace to the start of compression, the specimens were overheated to 4 degrees above the intended testing temperature. As a result, the guaranteed temperature for the beginning of deformation was  $400\pm 2$  and  $500\pm 2$  °C. Graphite was used as a lubricant between the specimen and the heads, with the Coulomb friction coefficient  $\mu$ =0.1 at temperatures ranging between 300 and 400 °C and  $\mu$ =0.15 at 500 °C for aluminum alloys.





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The dotted curves in Fig. 1 show experimental data on compression. While constructing the graphical dependences, we consider that the experiment deals with a uniaxial stress state. The specimen strain rate  $\dot{\epsilon}$  is calculated as

$$\dot{\varepsilon}(t) = \dot{h}(t)/h(t), \qquad (1)$$

where h(t) is the current height of the specimen at a time moment t,  $h(t) = h_0 - \Delta h(t)$ . Strain resistance  $\sigma_s$  is calculated by the formula

$$\sigma_s(t) = P(t)/F(t), \qquad (2)$$

where P(t) is the current pressing force; F(t) is the current specimen cross-sectional area determined from the constant volume condition for the specimen and the assumption of its strain homo-

geneity under uniaxial stress,  $F(t) = \frac{\pi d_0^2 h_0}{4h(t)}$  ( $h_0$  is the initial height of the specimen;  $d_0$  is the ini-

tial diameter of the specimen). It is obvious from Fig. 1a that the strain rate of the specimens was varied during the experiment. The necessity of compressing specimens with variable strain rate within one experiment is a peculiarity of the model parameter identification procedure based on the principles of dynamic object identification. According to these principles, only one experiment is in principle enough to evaluate the parameters of a mathematical model, see [46]. The obtained parameter values are valid only in the investigated interval of strain rates. To extend this interval, experiments for each temperature were carried out with realizing two different functions of time-varying strain rate (hereinafter referred to as loading form) presented in Fig. 1a, at least four experiments being made for each form.

The microstructure was studied by the method of electron backscattered diffraction (EBSD) with the use of a MIRA 3 TESCAN raster electron microscope with an OXFORD HKLNordlysF+ detachable device installed at the Institute of Engineering Science, Russian Academy of Sciences (Ural Branch), and by the method of transmission electron microscopy (TEM) with the use of a JEM200CX microscope installed at the Testing Center for Nanotechnologies and Advanced Materials, a collective use center belonging to the Institute of the Physics of Metals, Russian Academy of Sciences (Ural Branch). The scanning step for the EBSD analysis was 300 nm. The microstructure analysis was made by one representative specimen in a series of similar experiments. We assumed that the grains had misorientation exceeding 15° and that the subgrain misorientation ranged from 2 to 15°. When restoring the specimen microstructure, we assumed that a grain could not be smaller than the double step of scanning, i.e. each grain must have at least two points with a certain crystallographic direction. The volume portion of dynamic recrystallization that had occurred was calculated with the use of data on the mean misorientation of the lattice inside a grain  $\Theta_{mid}$  and the critical misorientation angle  $\Theta_{max}$  [47]. In [45] a technique for calculating  $\Theta_{max}$  was proposed, and its value for the AMg6 alloy was found to be 1.3°. To determine the value of  $\Theta_{max}$ , the middle misorientation angle  $\Theta_{mid}$  was calculated for each grain. Then, using a search with a step of 0.05, we calculated the critical misorientation angle  $\Theta_{max}$  so that the increase or decrease in the value of  $\Theta_{max}$ does not turn the deformed grains into recrystallized grains, and vice-versa. Figure 2 shows an example of determining recrystallized and deformed grains. The portion of dynamically recrystallized grains  $V_{exp}$  is here calculated as the ratio between the total area of recrystallized grains and the total area of grains.

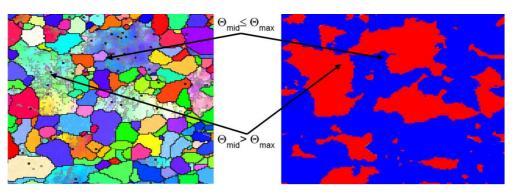


Fig. 2. An example of determining the recrystallized (*blue*) and deformed (*red*) grains when  $\Theta_{max}=1.3^{\circ}$ 

## 3. Results and discussion

## 3.1 Analysis of hardening and softening

The specimen microstructure was analyzed in the specimen center. Figure 3 shows the microstructure before and after deformation of the specimens. On the microstructure images in Fig. 3, the pixel color corresponds to definite crystallographic orientation. Each crystallographic orientation is represented in the color scale of the inverse pole figure, see Fig. 3*d*.

Figure 3*a* shows the microstructure of the AMg6 alloy specimen before deformation. It is obvious from the microstructure images that in the specimen center there are large elongated grains. Figs 3*b* and 3*c* demonstrate the microstructure images in the center of the specimens compressed according to form II (Fig. 1) at temperatures of 400 and 500 °C, respectively. Figure 4 shows a distribution map of deformed and recrystallized grains at  $\Theta_{max}$ =1.3° after specimen compression. In Figs 3 and 4 big black spots symbolize areas with indefinite crystallographic orientation and, generally, they are dispersoids. According to [48], the dispersoids may be the following compounds: Al<sub>3</sub>Mg<sub>2</sub>, Al<sub>6</sub>Mn, Mg<sub>2</sub>Si and AlFeSiMn.

Figure 4 shows that the portion represented by recrystallized grains in the residual microstructure of the deformed specimens is significant. Thus the volume portion of dynamic recrystallization  $V_{exp}$  is 82 % for 400 °C and 90 % for 500 °C. The portion of dynamically recrystallized grains is here calculated as the ratio between the total area of recrystallized grains and the total area of grains.



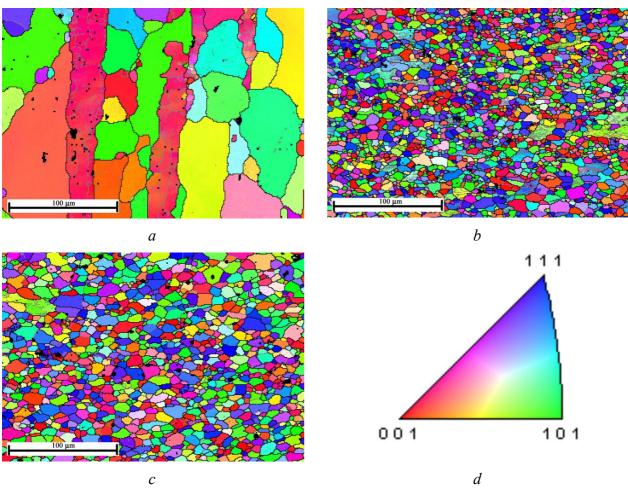
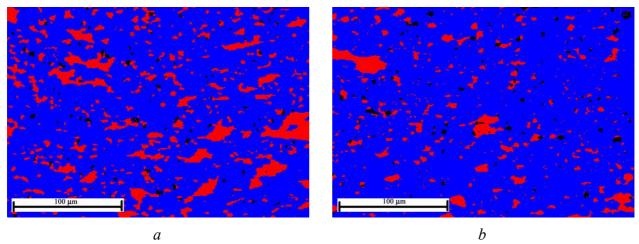
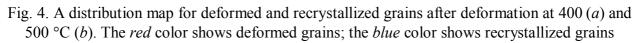


Fig. 3. The microstructure of an AMg6 alloy specimen in the central zone before deformation (a) and after deformation according to loading form II (Fig. 1*a*) at 400 (*b*) and 500 °C (*c*). The inverse pole figure (*d*)





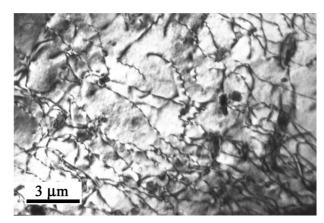


Fig. 5. Dispersoids and dislocations in the AMg6 alloy

The presence of a large fraction of recrystallized grains in the residual microstructure suggests that, at high strain, the growth of the grains was delayed by the growth of the neighboring grains [1–3]. As a result, the rate of material softening gradually decreased and that may cause material hardening. However, material hardening may be induced not only by decelerating recrystallization, but also by the presence of a large number of inclusions in the specimen microstructure. These inclusions block the motion of free dislocations and prevent the motion of boundaries and subboundaries. In the specimen microstructure before and after deformation, there are dispersoids sized between 60 nm and 2  $\mu$ m, see Fig. 5. According to the images made by a transmission electron microstructure during deformation. Therefore all the dispersoids in the alloy were formed in the previous stages of material treatment. However, the large density of dispersoids gives rise to a significant number of areas where dislocations, which were involved in plastic deformation, were blocked and that is certain to have resulted in the hardening of the material.

#### 3.2 Strain resistance model

The paper [42] proposed a strain resistance model describing the rheological behavior of alloys in conditions when an alloy exhibited viscous properties during plastic deformation, when hardening took place due to increasing dislocation density and softening occured due to dynamic recrystallization and recovery. The works [44, 45 and 49] showed that the developed model described, with a good accuracy, the rheology of the AMg6 alloy at a temperature of 300 °C in a wide strain rate range. It was found that the main process of softening was dynamic recrystallization, which occurred to a significantly smaller degree than at temperatures of 400 and 500 °C. At strain temperatures of 400 and 500 °C, the recrystallization was active, and the influence of the barrier effect on the growth of grains and the blocking of free dislocations increased. As a result, the absence of structural constituents in the strain resistance model, which would describe the barrier effect, caused the incorrect description of the strain resistance curve.

A new strain resistance model has been made on the basis of the previously developed model described in [45]. The entire volume of an alloy is divided into three portions named  $V_r$ ,  $V_p$  and  $V_n$ . The value  $V_r$  is the volume portion with dynamic recrystallization; the value  $V_p$  is the volume portion with dislocations blocked by dispersoids and impurity atoms; the value  $V_n$  is equal to the remaining portion of the alloy volume. This decomposition satisfies the equality

$$V_n + V_r + V_p = 1. (3)$$

At the initial moment of time before deformation,



It is assumed that all the plastic deformation is concentrated in the volume  $V_n$ , that the increment of yield stress due to the shear k in the Mises plasticity condition is attributed to plastic deformation and that this increment is proportional to the volume  $V_n$ , the increment of dislocation density and the volume  $V_p$ .

During plastic deformation, free dislocations involved in plastic deformation are pinned by dispersoids or impurity atoms thus causing material hardening [4, 13]. At the same time, at high temperatures due to high mobility of dislocations, dynamic recovery occurs actively, and this results in the annihilation of pinned dislocations and, as a consequence, in the reduction of the volume  $V_p$ . Thus, during plastic deformation, the volume  $V_p$  can both increase and decrease due to two competing processes.

Based on the above, the strain resistance model taking into account the barrier effects is presented in the following form:

$$\sigma_{s} = \sqrt{3}k + \frac{\sqrt{3}}{2}q;$$

$$k = a_{0} + \rho V_{n}(a_{10} + a_{11}V_{\rho})^{a_{12}};$$

$$\dot{\rho} = a_{1} \exp(-a_{2}\rho)\dot{\varepsilon} - a_{3}\rho;$$

$$\dot{V}_{p} = a_{13}\dot{\varepsilon} - a_{14}\frac{V_{n}}{1+\rho};$$

$$\left\{ \begin{aligned} a_{5}\varepsilon_{r}R^{2}\frac{dR}{dt} & \text{if } V_{r} \leq a_{6}; \\ V_{r}^{*} & \left\{ \frac{1-V_{r}}{1-a_{6}} \right\}^{a_{7}} & \text{if } V_{r} > a_{6}; \\ V_{r}^{*} & \text{when } V_{r} = a_{6}; \\ \dot{V}_{r}^{*} & \text{when } V_{r} = a_{6}; \\ \varepsilon_{r} = \int_{0}^{t_{r}} \dot{\varepsilon}dt; & \frac{dR}{dt} = \dot{\varepsilon}\rho \quad \text{when} \quad \rho > a_{4}; \\ q = a_{8} \ln(1+a_{9}\dot{\varepsilon}). \end{aligned} \right\}$$
(5)

Here *q* is a function describing the viscous properties of the material;  $\rho$  is the quantity proportional to the dislocation density increment induced by plastic deformation;  $\dot{\varepsilon}$  is strain rate (for the uniaxial stress state under compression  $\dot{\varepsilon} = |v|/h$ ; *v* is the speed of the testing machine grip; *h* is the current height of the specimen under deformation);  $\varepsilon_r$  is the strain accumulated before the onset of dynamic recrystallization; *R* is the radius of a recrystallized grain,  $R(t_r) = 0$ ,  $t_r$  is the time moment of the dynamic recrystallization onset defined by the condition  $\rho = a_4$ ;  $a_i$  (i = 0, ..., 14) are the model parameters to be identified by the experimental data. The dot above the symbols indicates a time derivative.

The system of equations (5), together with conditions (3) and (4), represents the physicalmechanical hierarchical rheological model describing, in the aggregate, the viscous and plastic properties of a medium. The model deals with plastic hardening due to the dislocation density increment and the blocking of the motion of free dislocations by dispersoids, as well as softening due to dynamic recovery and recrystallization.

# 3.3 Identification of the strain resistance model

The model parameters have been found by minimizing the root-mean-square deviation of the calculated values of strain resistance  $\sigma_s(t)$  from the experimental ones z(t) simultaneously with respect to two experimental strain resistance curves for each temperature separately, as

$$J(a_0, \dots, a_{14}) = \sum_{j=1}^2 \int_0^T \left[ \sigma_s(t) - z(t) \right]^2 dt, \qquad (6)$$

where T is the time of specimen deformation. The identification results for the strain resistance model are shown in Fig. 6b by solid lines. The model parameters are given in the Table.

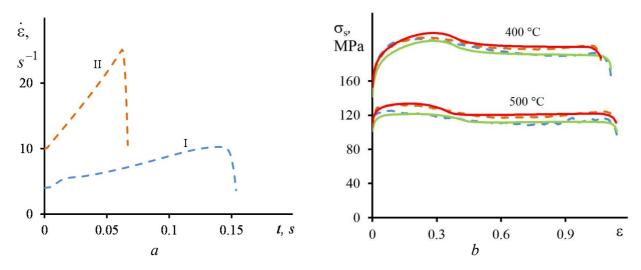
<i>T</i> , °C	$a_0$	$a_1$	$a_2$	$a_3$	$a_4$	$a_5$	$a_6$	$a_7$	$a_8$	<i>a</i> 9	$a_{10}$	$a_{11}$	<i>a</i> <sub>12</sub>	<i>a</i> <sub>13</sub>	<i>a</i> <sub>14</sub>
400	64	10125	0.03	0.003	8.2	0.009	0.254	4.58	13.8	3.6	0.3	1.4	1.13	0.04	0.14
500	44	10577	0.02	86.747	14.7	0.021	0.67	2.52	13.4	1.9	0.24	0.2	1.62	0.04	0.27

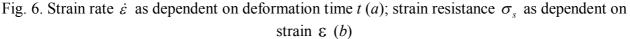
Table - The parameters of the strain resistance model

The average relative deviation  $\delta$  of the experimental curves from the calculated ones is 1.4 % for the test temperature of 400 °C and 1.6 % for the test temperature of 500 °C. The value of  $\delta$  is calculated by the formula:

$$\delta = \frac{1}{N} \sum_{j=1}^{N-2} \frac{\int_{0}^{T_{j}} |\sigma_{s}(t) - z(t)| dt}{\int_{0}^{T_{j}} z(t) dt} \cdot 100\%,$$
(7)

where j is the experiment number; N is the total number of tests in all the experiments used for identification (N = 2);  $T_j$  is the time of specimen deformation for the j-th experiment; z denotes the experimental values of strain resistance and  $\sigma_s$  denotes the values obtained from the identification of the strain resistance model for the j-th experiment.







The experimental data is indicated by the dotted curve and the calculated one is shown by the solid curve. The blue dotted curve corresponds to loading form I and the orange one corresponds to loading form II.

Figure 6b shows that strain resistance increases at the beginning, then reaches its peak value, decreases and increases again. The proposed mathematical model of strain resistance takes into account this rheological behavior of the material. This means that the model reflects qualitatively correctly the physical processes of structure formation in the alloy during deformation.

The obtained parameters of the strain resistance model (see Table) can be used to construct strain resistance curves for a constant strain rate within the range of strain rates used in the experiment. Figure 7 shows strain dependences of strain resistance  $\sigma_s$  at constant strain rates equal to 5, 10 and 20 s<sup>-1</sup>.

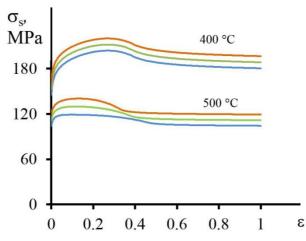


Fig. 7. Strain resistance  $\sigma_s$  as dependent on strain  $\varepsilon$  at a constant strain rate, s<sup>-1</sup>: 5 (*blue* curve), 10 (*green* curve) and 20 (*orange* curve)

#### 4. Conclusion

1. In this paper we have constructed a strain resistance model describing the rheological behavior of alloys at high-temperature plastic deformation when the processes of hardening due to the dislocation density increment, dynamic recovery, recrystallization and the blocking of the motion of free dislocations by inclusions occur simultaneously in a material.

2. It is illustrated by the AMg6 alloy that the strain resistance model qualitatively and quantitatively correctly describes the behavior for the strain resistance curve of the alloy at temperatures of 400 and 500 °C when  $\dot{\varepsilon} = 5 - 25 \text{ s}^{-1}$ .

3. Applying the method of electron backscattered diffraction, we have established that the main process of AMg6 alloy softening at temperatures of 400 and 500°C when  $\dot{\varepsilon} = 5 - 25 \text{ s}^{-1}$  is dynamic recrystallization resulting in the formation of new grains. The volume fraction of dynamic recrystallization is 80 and 90 % of the total volume at test temperatures of 400 and 500 °C, respectively.

4. The appearance of the second hardening segment on the strain resistance curve (see Fig. 1*b*) is caused by the slowdown of dynamic recrystallization and the barrier effect of dispersoids blocking free dislocations as well as grain and subgrain boundaries.

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